

# Technical Notes

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## Performance of a Low-Power Cylindrical Hall Thruster

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### Nomenclature

$B$ ,  $B$  = magnetic induction, T  
 $I$  = discharge current, A  
 $I_{sp}$  = specific impulse, s  
 $\dot{m}$  = flow rate, sccm  
 $P$  = discharge power, W  
 $V$  = discharge voltage, V  
 $\mu_e$  = magnetic moment, J/T

### I. Introduction

RECENT mission studies [1,2] have shown that a Hall thruster that operates at relatively constant thrust efficiency (45–55%) over a broad power range (300 W–3 kW) is enabling for deep space science missions when compared with state-of-the-art ion thrusters. Although conventional (annular) Hall thrusters can operate at high thrust efficiency at kilowatt power levels, it is difficult to construct one that operates over a broad power envelope down to  $\mathcal{O}(100\text{ W})$  while maintaining relatively high efficiency [3]. Scaling to low power requires a decrease in the thruster channel size and an increase

in the magnetic field strength while holding the dimensionless performance scaling parameters constant [4,5]. Increasing the magnetic field becomes technically challenging because the field can more easily saturate the miniaturized inner components of the magnetic circuit, and scaling down the magnetic circuit leaves very little room for magnetic pole pieces and heat shields. This makes it difficult to arrive at an optimal magnetic field configuration. Nonoptimal fields lead to enhanced power and ion losses that lower efficiency and result in increased heating and erosion of thruster components, particularly the critical inner components comprising the coaxial channel and magnetic circuit. Erosion of the thruster channel is one of the main life-limiting factors in conventional Hall thrusters [4].

An alternative approach to the miniaturization problem that has been demonstrated is embodied in the cylindrical-geometry Hall thruster (CHT) [6], which has recently been further developed especially for low-power applications [7]. Figure 1a shows that the CHT departs from the conventional, purely annular, Hall thruster geometry [4]. In contrast to the conventional annular geometry, in the cylindrical geometry the axial potential distribution is critical for electron confinement. This is because there is a large axial gradient in the magnetic field over the cylindrical part of the channel. The electrons drift both azimuthally around the thruster axis and outwards through the  $\mu_e \nabla B$  force. In the absence of an axial potential, the electrons would mirror out of the region of high magnetic field. In the CHT, the axial potential that accelerates ions outwards also plays an important role in confining the electrons within the thruster by counteracting the mirroring effect. In addition to being unlike an annular Hall thruster, the CHT differs from an end Hall thruster [8], which has a purely cylindrical geometry, biased channel walls, and a mostly axial applied magnetic field. This is in contrast to the cylindrical ceramic channel with a short annular region to sustain ionization and produce a magnetic field with a strong radial component in the CHT.

The detailed physics in the discharge channel of the CHT has been extensively studied [6,7,9]. In this Note, we report the measured performance ( $I_{sp}$ , thrust, and efficiency) of a cylindrical Hall thruster operating at  $\mathcal{O}(100\text{ W})$  input power.

### II. Experimental Apparatus

Performance measurements are obtained for a 3-cm Princeton Plasma Physics Laboratory (PPPL) CHT, shown in Fig. 1b. The thruster consists of a boron-nitride ceramic channel, an annular anode, two electromagnet coils, and a magnetic core. The thruster channel is a composite of a shorter annular region and a longer cylindrical region. Gas is injected through the anode into the short annular region of the thruster. The length of the annular region is sized to be greater than the ionization mean free path for xenon. This allows for high ionization of the propellant at the boundary between the annular and cylindrical regions. The electromagnet coils are operated using independently controlled power supplies, and the resulting field topology has a mirrorlike structure near the thruster axis.

The working propellant for these experiments is research-grade xenon gas. The cathode and anode flow rates are independently controlled using two variable 10-sccm MKS flow controllers (calibrated on Xe and controllable to  $\pm 0.1$  sccm). A commercial HeatWave Labs HWPES-250 hollow cathode is used in these

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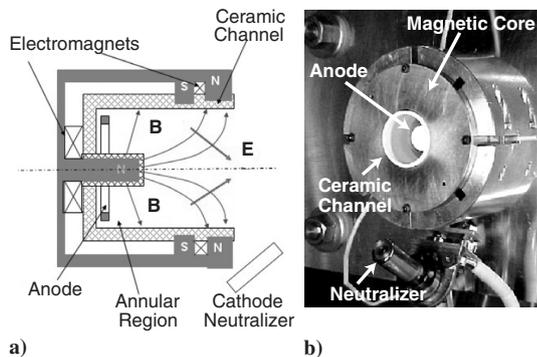


Fig. 1 PPPL cylindrical Hall thruster: a) schematic illustration and b) photograph.

experiments, serving as both the thruster cathode and the neutralizer. In all experiments, the cathode flow rate is 1 sccm. The cathode placement and operation was quite similar to that used in previous testing of the 2.6-cm PPPL CHT [7,9]. The main difference, which did not produce any effect on the discharge characteristics in our testing, is that in the previous experiments the cathode flow rate was typically 2 sccm.

Thrust was measured using the variable-amplitude hanging pendulum with extended-range (VAHPER) thrust stand [10] at the NASA Marshall Space Flight Center (MSFC). This thrust stand employs a unique mechanical linkage system that converts horizontal deflection of the pendulum arm into amplified vertical deflection. The thrust stand is mounted inside a 9-ft-diam, 25-ft-long stainless steel vacuum chamber. The vacuum level inside the chamber is maintained by two 2400 l/s turbopumps and two 9500 l/s cryopumps. The base pressure was  $5.7 \times 10^{-7}$  torr and the background pressure of xenon during thruster operation at  $\sim 5$  sccm total flow rate was  $\sim 9 \times 10^{-6}$ .

### III. Experimental Results and Discussion

Displacement (thrust) calibration of the VAHPER thrust stand is accomplished using an in situ calibration rig that applies a series of known loads normal to the pendulum arm. Calibration can be performed before, during, and after thruster operation. The measured displacement of the vertically deflecting linkage is recorded as the calibration loads are applied to the arm. Assuming that the relationship between the applied force and the measured displacement is linear allows for a linear curve fit of the calibration data found in Fig. 2. The calibration routine propagates the errors associated with the individual force and displacement measurements to accurately account for all systematic errors [11].

Thrust is found by first calculating the linear gap displacement transducer (LGDT) output voltage difference between the steady-state portion of thruster operation and the zero level output immediately following thruster shutdown (see Fig. 3). These data are then converted to thrust using the calibration curve-fit constants found in Fig. 2. Anode efficiency and  $I_{sp}$  (specifically, anode  $I_{sp}$ ) are

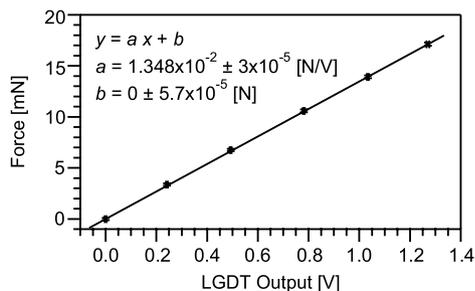


Fig. 2 Applied calibration force plotted as a function of LGDT response. The linear curve is a fit to the displayed data. Error bars on the data points and the fit coefficients represent a 68% ( $1\sigma$ ) confidence interval.

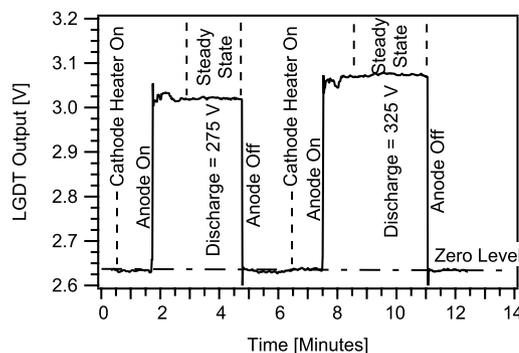


Fig. 3 Raw LGDT output data as a function of elapsed time for two thruster operating points.

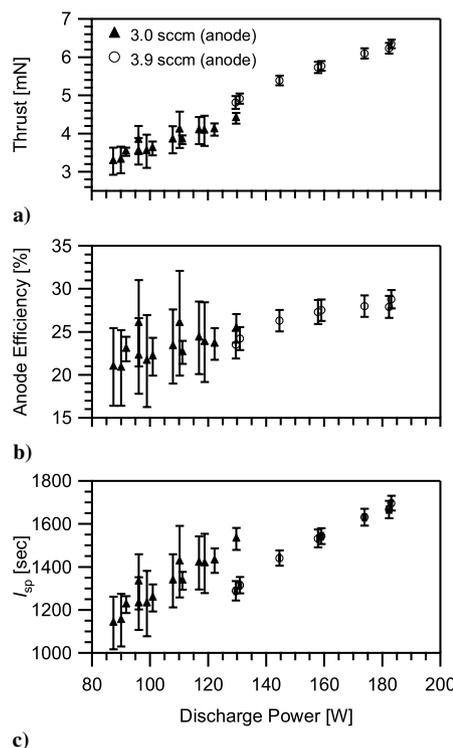


Fig. 4 Performance measurements for a 3-cm cylindrical Hall thruster: a) thrust, b) anode efficiency, and c)  $I_{sp}$  as a function of discharge power. The error bars represent a 95% ( $2\sigma$ ) confidence interval.

computed according to their standard definitions [12] using recorded anode mass flow rates and power supply outputs. The electromagnet power is not included in the anode efficiency calculation. The uncertainty levels on the performance data represent a 95% ( $2\sigma$ ) confidence interval.

Two sets of thrust measurements corresponding to two anode mass flow rates (3 and 3.9 sccm) are presented in Fig. 4a, with the corresponding anode efficiencies plotted in Fig. 4b and  $I_{sp}$  values

Table 1 Nominal thruster operating conditions for each of the anode flow rates tested

$\dot{m} = 3$ sccm			$\dot{m} = 3.9$ sccm		
V, V	I, A	P, W	V, V	I, A	P, W
250	0.36	90	250	0.52	130
275	0.36	99	275	0.53	146
300	0.37	111	300	0.53	159
325	0.38	123	325	0.54	176
340	0.38	129	340	0.54	184

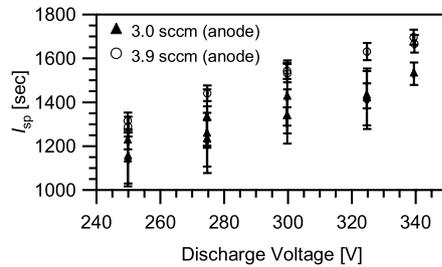


Fig. 5  $I_{sp}$  as a function of discharge voltage. The error bars represent a 95% ( $2\sigma$ ) confidence interval.

found in Fig. 4c. These data span a range between 90 to 185 W in discharge power and result in thrust levels between 3 and 6 mN, anode efficiencies between 20 and 27%, and  $I_{sp}$  levels between 1100 and 1650 s. For completeness, the voltage, current, and power levels associated with these data are given in Table 1. At the power levels tested, the 3-cm PPPL CHT performance compares favorably with both the SPT-30 [13] and BHT-200 [14] Hall thrusters and is significantly greater than the 6% efficiency and  $I_{sp}$  of 850 s measured in a miniaturized conventional (coaxial) Hall thruster at 125-W input power [5]. It is also comparable to previous performance measurements of the 2.6-cm PPPL CHT [7].

In Fig. 4a, thrust increases with discharge power. Anode efficiency (Fig. 4b) also increases with discharge power, going from 20 to 27%. In addition, it appears to asymptote at the higher power levels. Specific impulse (Fig. 4c) increases with discharge power. Additional insight is provided in Fig. 5, in which we see that for a constant mass flow rate,  $I_{sp}$  increases with discharge voltage. However, for a given voltage,  $I_{sp}$  is greater at the higher mass flow rate. We can speculate that  $I_{sp}$  increases due to greater propellant use (i.e., low neutral fraction), the formation of doubly ionized Xe, or a combination of the two. However, we do not have enough information at this time to reach a definite conclusion.

In Fig. 4, we observe that the lower-power data points generally possess much larger errors than the higher-power points. The larger errors are partially attributable to recording of insufficient significant figures in the discharge current measurement for those operating points. The issue was resolved before higher-power testing, but the error had to be conservatively overestimated for the lower-power data set. Additional errors in this lower-power set are due to thermally induced deflections of the vacuum chamber, which occurred as the interior of the building was heated by the sun and warm ambient air and subsequently cooled by the HVAC system, which blew directly onto the vacuum chamber. Both heating and cooling effects were noticeable within minutes and appeared as drifts in the thrust stand position. The heating effects were reduced during the second, higher-power, trial by testing at night and adjusting the HVAC ducts so that they did not blow directly onto the vacuum chamber.

#### IV. Conclusions

Although conventional (annular) Hall thrusters are efficient in the kilowatt power regime, they become inefficient when scaled down to small sizes. This is due to the difficulties associated with holding the performance scaling parameters constant while decreasing the channel size and increasing the applied magnetic field strength. The cylindrical Hall thruster can be more readily scaled to smaller sizes due to its nonconventional discharge-chamber geometry and associated magnetic field profile.

A series of performance measurements on the 3-cm PPPL CHT over a power ranging from ~85–185 W were performed using the

VAHPER thrust stand. The PPPL CHT produced thrust levels ranging from 3–6 mN, anode efficiencies spanning 20–27%, and  $I_{sp}$  between 1100–1650 s. Thrust increased as a function of discharge power, whereas  $I_{sp}$  increased with discharge voltage and with increasing anode flow rate.

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